

# On the Subcritical SBBM (Self-Catalytic Branching Brownian Motions)

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*Joint ongoing work with Haojie Hou*

- **Model: Self-Catalytic Branching Brownian Motions (SBBM)**

A particle system that extends the *Branching Brownian Motion* by incorporating *catalytic branching* through pairwise interactions between particles (which will be defined in the next slide).

- **Objectives:**

- Construct an SBBM that supports *infinitely many particles*.
- Characterize the *law* of the SBBM.
- Investigate the *coming down from infinity (CDI)* property of the SBBM.

- **Motivation:**

- SBBMs offer insights into complex population dynamics, incorporating *biological dispersal* and *intraspecific competition/cooperation*.
- SBBMs serve as moment duals to a class of *stochastic reaction-diffusion equations with multiplicative noise*, providing valuable information on the *well-posedness*, *propagation speed*, and *compact support properties* of SPDEs.

## SBBM Dynamics:

- **Initial Configuration:**

The positions of the initial particles are given by  $(x_i)_{i=1}^n \subset \mathbb{R}$ .

- **Particle Movement:**

Each particle performs independent Brownian motion on  $\mathbb{R}$ .

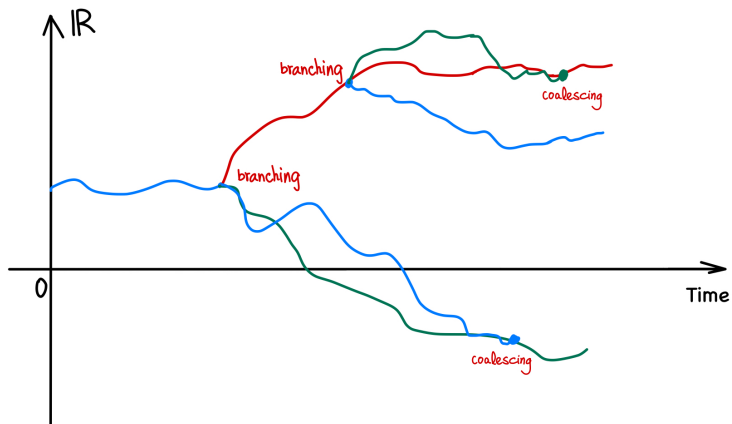
- **Ordinary Branching:**

Particles branch independently at rate  $\beta_o$ , replaced by  $k$  offspring according to the law  $(p_k)$ .

- **Catalytic Branching:**

Each pair of particles independently branches at rate  $\beta_c$ , based on their intersection local times, replaced by  $k$  offspring according to the law  $(q_k)$ .

# An illustration of SBBM



- A graph illustration of SBBM with  $p_3 = 1$  and  $q_1 = 1$ .

- **Catalytic Branching is Subcritical:**

$$\sum kq_k < 2.$$

- **Catalytic Branching is Not Parity-Preserving:**

There exists an odd  $k$  such that  $q_k > 0$ .

- **A Technical Assumption:**

There exists  $R > 1$  such that

$$\sum R^k p_k < \infty \quad \text{and} \quad \sum R^k q_k < \infty.$$

# The Explosion Problem for SBBM

## A Priori Consideration:

- The SBBM model is well-defined only up to its explosion time  $\tau_\infty$ .

## Definition of the Explosion Time $\tau_\infty$ :

- Define a sequence of stopping times  $\{\tau_k\}_{k \in \mathbb{Z}_+}$  such that:
  - $\tau_0 = 0$ .
  - For each  $k \geq 0$ ,

$$\tau_{k+1} := \inf\{t > \tau_k : \text{a branching occurs at time } t\}.$$

- The explosion time is defined as:

$$\tau_\infty := \lim_{k \rightarrow \infty} \tau_k.$$

## The Explosion Problem:

- Does  $\tau_\infty = \infty$ ?

# Non-Explosion Property

## Athreya-Tribe (2000, Ann. Probab.)

If the ordinary branching is subcritical, i.e.,  $\sum kp_k < 1$ , then:

$$\mathbb{P}(\tau_\infty = \infty) = 1.$$

## Hou-Sun (ongoing)

The result above holds without assuming  $\sum kp_k < 1$ .

### Implications:

- The SBBM model remains well-defined for all time.
- For  $t \geq 0$ , define  $Z_t^{(n)}(A) := \#\{\text{particles in set } A \text{ at time } t\}$ .
- The process  $(Z_t^{(n)})_{t \geq 0}$  is a Markov process taking values in:

$$\mathcal{N} := \{\text{locally finite point measures on } \mathbb{R}\}.$$

- We refer to  $Z_t^{(n)}$  as an SBBM (with  $n$  initial particles).

# The Infinitely Initial Particles Problem

- Consider the scenario where the initial configuration of an SBBM consists of infinitely many particles located at  $(x_i)_{i=1}^{\infty} \subset \mathbb{R}$ .
- In this case,  $\tau_{\infty} = 0$  almost surely, and the model is a priori not well-defined.

## The Infinitely Initial Particles Problem:

- Let  $(Z_t^{(n)})_{t \geq 0}$  be an SBBM with initial value  $\sum_{i=1}^n \delta_{x_i}$ .
- What is the limit of the processes  $(Z_t^{(n)})_{t \geq 0}$  as  $n \rightarrow \infty$ ?

## Characterizing the Law:

- How can we characterize the law of  $(Z_t^{(n)})_{t \geq 0}$ ?
- Does the processes  $\{(Z_t^{(n)})_{t \geq 0} : n \in \mathbb{N}\}$  satisfy the tightness property?



# The Dual SPDE

- **The dual SPDE of SBBM:**

$$\begin{cases} \partial_t u_t(x) = \frac{\Delta}{2} u_t(x) - \Phi(u_t(x)) + \sqrt{\Psi(u_t(x))} \dot{W}_t(x), & t > 0, x \in \mathbb{R}, \\ u_0(x) = f(x), & x \in \mathbb{R}. \end{cases}$$

- **Space-time White Noise:**

$(W_t)_{t \geq 0} :=$  a cylindrical Wiener process on  $L^2(\mathbb{R})$ , such that

$$\mathbb{E}[W_t(\phi)W_s(\psi)] = (t \wedge s)\langle \phi, \psi \rangle.$$

- **Ordinary Branching Mechanism:**

$$\Phi(z) := \beta_o \left( \sum_{k=0}^{\infty} p_k (1-z)^k - (1-z) \right).$$

- **Catalytic Branching Mechanism:**

$$\Psi(z) := \beta_c \left( \sum_{k=0}^{\infty} q_k (1-z)^k - (1-z)^2 \right).$$

# The Dual SPDE

- Let  $z^* := \inf\{z \in [1, 2] : \Psi(z) = 0\}$ .
- Under the assumption that the catalytic branching is not parity-preserving, we have  $z^* \in [1, 2)$ .
- Let  $C(\mathbb{R}, [0, z^*]) := \{\text{Continuous functions from } \mathbb{R} \text{ to } [0, z^*]\}$ .
- We say a  $C(\mathbb{R}, [0, z^*])$ -valued continuous process  $(u_t)_{t \geq 0}$  is a weak solution to the dual SPDE if there exists a space-time white noise  $W$  such that for all  $\phi \in C_c(\mathbb{R})$ , almost surely,

$$\begin{aligned} \langle u_t, \phi \rangle - \langle f, \phi \rangle &= \int_0^t \langle u_s, \frac{\Delta}{2} \phi \rangle ds - \int_0^t \langle \Phi(u_s), \phi \rangle ds \\ &\quad + \int_0^t \langle \sqrt{\Psi(u_s)} \phi, dW_s \rangle. \end{aligned}$$

**Shiga (1994, Can. J. Math.)**

For each initial value  $f \in C(\mathbb{R}, [0, z^*])$ , there exists a weak solution to the dual SPDE.

# The Duality

- For any  $[0, z^*]$ -valued function  $g$  and point measure  $\mu$ , define

$$(1 - g)^\mu := \prod_{x \in \mathbb{R}} (1 - g(x))^{\mu(\{x\})}.$$

## Athreya-Tribe (2000, Ann. Probab.)

If the ordinary branching is subcritical, i.e.,  $\sum kp_k < 1$ , then:

$$\mathbb{E} \left[ (1 - u_t)^{Z_0^{(n)}} \right] = \mathbb{E} \left[ (1 - u_0)^{Z_t^{(n)}} \right].$$

## Hou-Sun (ongoing)

The result above holds without assuming  $\sum kp_k < 1$ .

- Corollary:** The uniqueness in law holds for the dual SPDE.

# SBBM with Infinitely Many Initial Particles

- Let the state space  $\mathcal{N}$  be equipped with the vague topology.

**Initial Trace**  $(\Lambda, \mu)$ :

- $\Lambda := \{\text{sub-sequential limits of } (x_i)_{i=1}^{\infty}\}$ .
- $\mu := \sum_{x_i \notin \Lambda} \delta_{x_i}$ .

## Hou-Sun (ongoing)

There exists an  $\mathcal{N}$ -valued càdlàg Markov process  $(Z_t)_{t>0}$  such that  $(Z_t^{(n)})_{t>0}$  converges to  $(Z_t)_{t>0}$  as  $n \rightarrow \infty$  in finite-dimensional distributions. The law of the process  $(Z_t)_{t>0}$  is determined by the two branching mechanisms  $(\Phi, \Psi)$  and the initial trace  $(\Lambda, \mu)$ .

- We call  $(Z_t)_{t>0}$  an SBBM with initial trace  $(\Lambda, \mu)$  and branching mechanisms  $(\Phi, \Psi)$ .

# Coming Down from Infinity (CDI)

## The Local-Time Coalescing Brownian Motions (LCBM):

- If the ordinary branching rate  $\beta_o = 0$  and the catalytic branching law satisfies  $q_1 = 1$ , then the SBBM degenerates into the LCBM.
- In this case,  $\Phi = 0$  and  $\Psi(z) = z(1 - z)$ .

### Barnes-Mytnik-Sun (2023, Ann. Probab.)

Suppose that  $(Z_t)_{t>0}$  is an LCBM. Let  $U$  be any open interval. Then, almost surely, for every  $t > 0$ ,

$$Z_t(U) < \infty \iff (\Lambda \cup \text{supp}(\mu)) \cap U \text{ is bounded.}$$

- **The CDI property:** Almost surely, for every  $t > 0$ ,

$$Z_t(\mathbb{R}) < \infty \iff \sup\{|x_i| : i \in \mathbb{N}\} < \infty.$$

### Hou-Sun (ongoing)

Same result holds for the SBBM.

# The Mean Field Equation (MFE)

From a physic's point of view:

- The MFE for a system of independent Brownian motions is given by the heat equation  $\partial_t h = \frac{\Delta}{2} h$ . In the sense that

$$\mathbb{E}[\#\{\text{particles in } (x - \frac{1}{2}, x + \frac{1}{2}) \text{ at time } t\}] \approx h_t(x).$$

- The MFE for LCBM is  $\partial_t v = \frac{\Delta}{2} v - \frac{\Psi'(0+)}{2} v^2$ .

Le Gall (1996, J. Appl. Math. Stochastic Anal.)

There exists a unique non-negative solution  $(v_t(x))_{t>0, x \in \mathbb{R}}$  to the PDE

$$\begin{cases} \partial_t v_t(x) = \frac{\Delta}{2} v_t(x) - \frac{\Psi'(0+)}{2} v_t(x)^2, & t > 0, x \in \mathbb{R}, \\ \left\{ y \in \mathbb{R} : \forall r > 0, \lim_{t \rightarrow 0} \int_{y-r}^{y+r} v_t(x) dx = \infty \right\} = \Lambda, \\ \lim_{t \rightarrow 0} \langle v_t, \phi \rangle = \langle \mu, \phi \rangle, & \phi \in C_c(\Lambda^c). \end{cases}$$

# Speed of CDI for LCBM

## The Speed of CDI Problem

- Assume CDI holds for a process  $(N_t)_{t \geq 0}$ .
- Can we find a rate function  $a(t)$  such that  $N_t/a(t) \rightarrow 1$  as  $t \downarrow 0$ ?

### Barnes-Mytnik-Sun (2023, Ann. Probab.)

Suppose that  $(Z_t)_{t > 0}$  is an LCBM with initial trace  $(\Lambda, \mu)$ . Let  $U$  be an open interval. Suppose that  $(\Lambda \cap \text{supp}(\mu)) \cap U$  is bounded and  $\Lambda \cap \bar{U} \neq \emptyset$ . Then,

$$\left( \int_U v_t(x) dx \right)^{-1} Z_t(U) \xrightarrow[t \downarrow 0]{L^1} 1,$$

where  $(v_t(x))_{t > 0, x \in \mathbb{R}}$  is the solution to the corresponding MFE with initial trace  $(\Lambda, \mu)$ .

### Hou-Sun (ongoing)

The exact same result holds for SBBM.

## Criticality of the Branching:

- It is crucial for our result that the catalytic branching is subcritical, i.e.,  $\sum kq_k < 2$ .
- [Barnes-Mytnik-Sun \(2024, Arxiv\)](#) constructed an SBBM with  $p_\infty = 1$  and  $q_1 = 1$ , and showed that the total population in this model is “reflecting from infinity”.
- When the catalytic branching is supercritical, i.e.,  $\sum kq_k > 2$ , we believe that the SBBM will explode in finite time.
- When there is no ordinary branching, i.e.,  $\beta_o = 0$ , and the catalytic branching is critical, i.e.,  $\sum kq_k = 2$ , we believe that the SBBM is non-explosive and rescales to the stochastic heat equation:

$$\partial_t u = \frac{\Delta}{2} u + u \dot{W}.$$



## About the Parity:

- It is crucial for our result that the catalytic branching is not parity-preserving, i.e., there exists an odd number  $k$  such that  $q_k > 0$ .
- Consider an SBBM with no ordinary branching, i.e.,  $\beta_0 = 0$ , and  $q_0 = 1$ . We call this model the local-time annihilating Brownian motion (LABM).
- LABM is non-explosive and can be defined up to all time, provided there are only finitely many initial particles.
- It can be shown that  $\{Z_t^{(n)} : n \in \mathbb{N}\}$  is tight.
- However, the subsequential convergence-in-distribution limit of  $\{Z_t^{(n)} : n \in \mathbb{N}\}$  is not unique.
- [Hammer-Ortgiese-Völlering \(2021, Stochastic Process. Appl.\)](#): The entrance laws of the (hard) annihilating Brownian motion are characterized.

# Comments

## Examples of Duality:

- We say two Markov processes  $(X_t)_{t \geq 0}$  and  $(Y_t)_{t \geq 0}$  are dual to each other if there exists a large class of functions  $H(x, y)$  such that

$$\mathbb{E}[H(X_t, Y_0)] = \mathbb{E}[H(X_0, Y_t)].$$

- [Bachelier \(1900, Ann. Sci. École Norm. Sup.\)](#): Brownian motion and the heat equation  $\partial_t h = \frac{\Delta}{2} h$ .
- [McKean \(1975, Comm. Pure Appl. Math.\)](#): Branching Brownian motion and the FKPP equation  $\partial_t v = \frac{\Delta}{2} v + v(1 - v)$ .
- [Harris \(1978, Ann. Probab.\)](#): Coalescing random walk and the voter model.
- [Shiga \(1986, Math. Appl.\)](#): LCBM and the stochastic FKPP equation  $\partial_t v = \frac{\Delta}{2} v + \sqrt{v(1 - v)} W$ .
- [Tóth-Werner \(1998, Probab. Theory Relat. Fields\)](#): (Hard) Coalescing Brownian motions and itself.
- **Folklore**: Stochastic heat equation  $\partial_t u = \frac{\Delta}{2} u + u \dot{W}$  and itself.
- ...

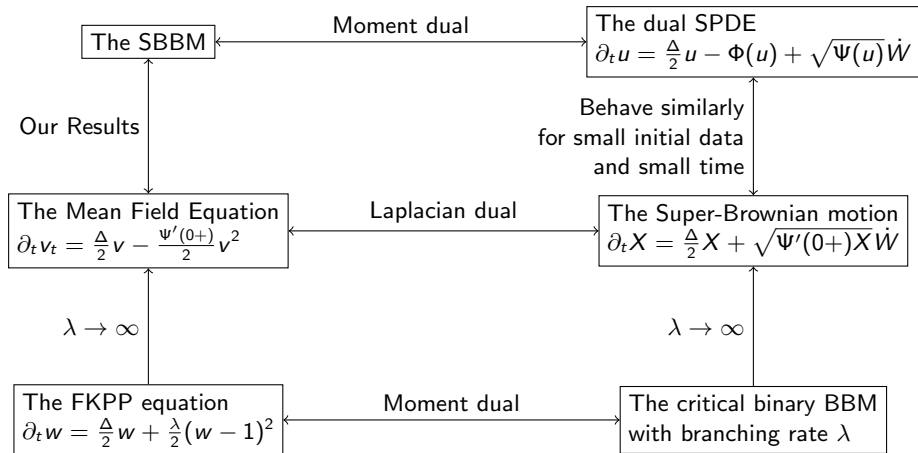
## Coming Down from Infinity (CDI):

- Aldous (1999, Bernoulli): Kingman's coalescent.
- Schweinsberg (2000, Electron. Comm. Probab.) and Berestycki-Berestycki-Limic (2010, Ann. Probab.):  $\Lambda$ -coalescent.
- Limic-Sturm (2006, Electron. J. Probab.) and Angel-Berestycki-Limic (2012, Probab. Theory Related Fields): Coalescing random walks on graphs.
- Murrat-Weber (2017, Comm. Math. Phys.): Dynamical  $\Phi_3^4$  model (leading to a new construction of the Euclidean  $\Phi_3^4$  Field Theory).
- Li-Yang-Zhou (2019, Ann. Appl. Probab.): Continuous-state nonlinear branching processes.
- Baguley-Döering-Shi (2024, Arxiv): Time-changed Lévy processes.
- ...

## The Mean Field Equation (MFE)

- The CDI rate of SBBM is characterized by  $\partial_t v = \frac{\Delta}{2} v - \frac{\Psi'(0+)}{2} v^2$  despite that the true MFE is  $\partial_t \tilde{v} = \frac{\Delta}{2} \tilde{v} + \Phi'(0+) \tilde{v} - \frac{\Psi'(0+)}{2} \tilde{v}^2$ .
- This is because  $v(s, y) \asymp \tilde{v}(s, y)$  uniformly for  $(s, y) \in [0, 1] \times \mathbb{R}$ .
- The equation  $\partial_t v = \frac{\Delta}{2} v - v|v|^\alpha$  with initial trace  $(\Lambda, \mu)$  was studied by [Marcus-Véron \(1999, Comm. Partial Differential Equations\)](#) in the PDE literature.
- [Watanabe \(1968, J. Math. Kyoto Univ.\)](#):  
The equation  $\partial_t v = \frac{\Delta}{2} v - \frac{\Psi'(0+)}{2} v^2$  is the Laplace dual to the Super-Brownian motion  $(X_t)_{t \geq 0}$ .
- [Le Gall \(1996, J. Appl. Math. Stochastic Anal.\)](#) used the equation  $\partial_t v = \frac{\Delta}{2} v - \frac{\Psi'(0+)}{2} v^2$  to study the Brownian snake, which is related to the super-Brownian motion through a Ray-Knight type theorem.

# Theory Roadmap



*Thanks!*